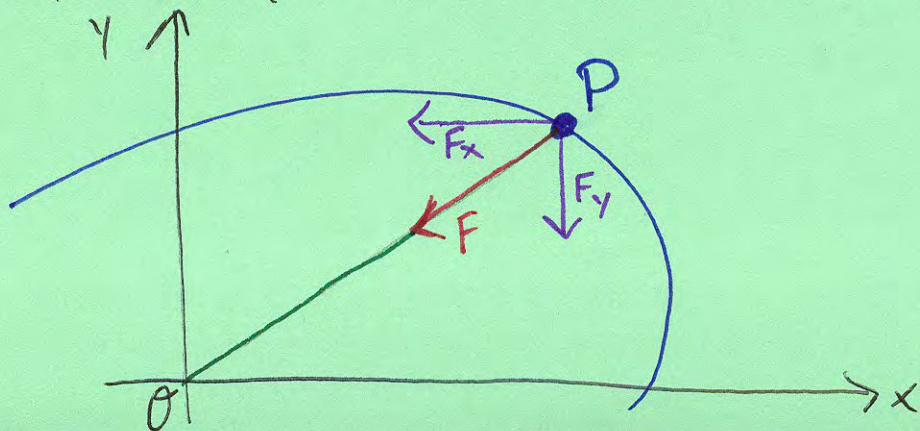


Centripetal Force in Cartesian Coordinates

Let P be a point-mass with mass m , moving in a plane, acted upon by a centripetal force with center the origin of the plane.

Let $x=x(t)$ and $y=y(t)$ be the coordinates of P as functions of time, $t \geq 0$. The magnitude of the force can also vary with time, so we write $F=F(t)$. We assume that the motion of P is smooth enough so that $x(t), y(t)$, are twice (at least) differentiable and that $F(t)$ is differentiable. Let r be the distance from O to P and let F_x and F_y be the components of F in the x & y directions, resp. Then r, F_x , and F_y are differentiable.



Using properties of similar triangles:

(2)

$$\frac{|F_x|}{F} = \frac{|x|}{r} \quad \& \quad \frac{|F_y|}{F} = \frac{|y|}{r}$$

Notice that the signs of F_x & x and F_y & y are always opposite, so that

$$\frac{F_x}{F} = -\frac{x}{r} \quad \& \quad \frac{F_y}{F} = -\frac{y}{r}$$

$$\Rightarrow rF_x = -xF \quad \& \quad rF_y = -yF$$

Now, Newton's second law gives:

$$F_x = m \frac{d^2x}{dt^2} \quad \& \quad F_y = m \frac{d^2y}{dt^2}$$

$$\text{Thus, } mr \frac{d^2x}{dt^2} = -xF \quad \& \quad mr \frac{d^2y}{dt^2} = -yF$$

Subtracting y times the first & x times the second:

$$mry \frac{d^2x}{dt^2} - mrx \frac{d^2y}{dt^2} = mr \left(y \frac{d^2x}{dt^2} - x \frac{d^2y}{dt^2} \right) = -yxF + -xyF = 0$$

(3)

This gives: $y \frac{d^2 x}{dt^2} = x \frac{d^2 y}{dt^2}$ (*)

We are interested in the quantity $x \frac{dy}{dt} - y \frac{dx}{dt}$ (it relates to $\frac{d\theta}{dt}$). Observe

$$\begin{aligned} \frac{d}{dt} \left(x \frac{dy}{dt} - y \frac{dx}{dt} \right) &= \cancel{\frac{dx}{dt} \frac{dy}{dt}} + x \frac{d^2 y}{dt^2} - \cancel{\frac{dy}{dt} \frac{dx}{dt}} - y \frac{d^2 x}{dt^2} \\ &= x \frac{d^2 y}{dt^2} - y \frac{d^2 x}{dt^2} = 0 \quad \text{by (*)} \end{aligned}$$

Therefore,

$$\boxed{x \frac{dy}{dt} - y \frac{dx}{dt} = c}$$

This relation tells us important information about the motion of P which we can decode by going to polar coordinates!

Let $r = f(\theta)$ be a function describing the motion of the particle P. Notice that

$\theta = \theta(t)$ & $r = r(t) = f(\theta(t))$. By our smoothness assumptions, $\theta, r,$ & f are twice differentiable.

We assume $t=0$ is an instant where P crosses the polar axis so that $\theta(0) = 0$. Take $\theta(t) > 0$ for $t > 0$ & $\theta(t) < 0$ for $t < 0$.

Now, let's change $x \frac{dy}{dt} - y \frac{dx}{dt} = c$ to polar:

$$\begin{aligned} x \frac{dy}{dt} &= r \cos \theta \frac{d}{dt}(r \sin \theta) = r \cos \theta \left(\frac{dr}{dt} \sin \theta + r \cos \theta \frac{d\theta}{dt} \right) \\ &= \cancel{\frac{dr}{dt} r \cos \theta \sin \theta} + r^2 \cos^2 \theta \frac{d\theta}{dt} \end{aligned}$$

$$\begin{aligned} y \frac{dx}{dt} &= r \sin \theta \frac{d}{dt}(r \cos \theta) = r \sin \theta \left(\frac{dr}{dt} \cos \theta - r \sin \theta \frac{d\theta}{dt} \right) \\ &= \cancel{\frac{dr}{dt} r \sin \theta \cos \theta} - r^2 \sin^2 \theta \frac{d\theta}{dt} \end{aligned}$$

$$x \frac{dy}{dt} - y \frac{dx}{dt} = r^2 \cos^2 \theta \frac{d\theta}{dt} + r^2 \sin^2 \theta \frac{d\theta}{dt} = r^2 \frac{d\theta}{dt}$$

So, we get

$$r^2 \frac{d\theta}{dt} = \boxed{(r(t))^2 \theta'(t) = c}$$

Now, consider P at two times t_0 & t_1 , and let A be the area swept out in this ^($t_0 \leq t_1$) time by OP . Let $\theta(t_0) = a$ & $\theta(t_1) = b$, so

$$A = \int_a^b \frac{1}{2} (f(\theta))^2 d\theta$$

Recall $r(t) = f(\theta(t))$. In terms of t :

$$d\theta = d(\theta(t)) = \theta'(t) dt$$

Thus

$$\begin{aligned} A &= \int_a^b \frac{1}{2} [f(\theta)]^2 d\theta = \int_{t_0}^{t_1} \frac{1}{2} r(t)^2 \theta'(t) dt \\ &= \int_{t_0}^{t_1} \frac{c}{2} dt = \frac{1}{2} c (t_1 - t_0) \end{aligned}$$

Let $k = \frac{1}{2} c$, then

$$\boxed{A = k(t_1 - t_0)}$$

So, the area swept out is equal to K times the time it takes to sweep it out!

Which law have we just verified?

Kepler's Second Law!

We call K the Kepler constant of the orbit.

Now we turn our focus to converting

$$mr \frac{d^2x}{dt^2} = -F_x \quad \& \quad mr \frac{d^2y}{dt^2} = -F_y \quad (**)$$

to polar.

$$\frac{d^2x}{dt^2} = \frac{d}{dt} \left(\frac{dr}{dt} \cos \theta - r \sin \theta \cdot \frac{d\theta}{dt} \right)$$

$$= \frac{d^2r}{dt^2} \cos \theta - \frac{dr}{dt} \sin \theta \frac{d\theta}{dt} - \frac{dr}{dt} \sin \theta \frac{d\theta}{dt} - r \cos \theta \left(\frac{d\theta}{dt} \right)^2 - r \sin \theta \frac{d^2\theta}{dt^2}$$

$$= \frac{d^2r}{dt^2} \cos \theta - r \cos \theta \left(\frac{d\theta}{dt} \right)^2 - \sin \theta \left(2 \frac{dr}{dt} \frac{d\theta}{dt} + r \frac{d^2\theta}{dt^2} \right)$$

Because $(r(t))^2 \theta'(t) = 2k$,

7

$$0 = \frac{d}{dt}(2k) = \frac{d}{dt} \left(r^2 \frac{d\theta}{dt} \right) = 2r \frac{dr}{dt} \frac{d\theta}{dt} + r^2 \frac{d^2\theta}{dt^2}$$

So, we have $\frac{d^2x}{dt^2} = \cos\theta \left[\frac{d^2r}{dt^2} - r \left(\frac{d\theta}{dt} \right)^2 \right]$

Similarly,

$$\frac{d^2y}{dt^2} = \sin\theta \left[\frac{d^2r}{dt^2} - r \left(\frac{d\theta}{dt} \right)^2 \right]$$

Putting these into ~~(*)~~ gives:

$$F \cos\theta = m \cos\theta \left[r \left(\frac{d\theta}{dt} \right)^2 - \frac{d^2r}{dt^2} \right] \quad \text{and}$$

$$F \sin\theta = m \sin\theta \left[r \left(\frac{d\theta}{dt} \right)^2 - \frac{d^2r}{dt^2} \right]$$

(after cancelling the r 's on both sides)

Since $\sin\theta$ & $\cos\theta$ are not simultaneously zero:

$$F = m \left[r \left(\frac{d\theta}{dt} \right)^2 - \frac{d^2r}{dt^2} \right]$$